

QUESTION

The wave equation for \vec{E} is

$$(\nabla^2 + k^2)\vec{E} = -i\omega\mu\vec{J}$$

which is the scalar equation $(\nabla^2 + k^2)u = f$ for each component.

So can we say that

$$\vec{E}(\vec{r}) = \iiint d\vec{r}' g(\vec{r}, \vec{r}') (-i\omega\mu\vec{J}(\vec{r}')) \quad ?$$

The answer is no. The reason is that the wave equation has more solutions than Maxwell's equations, e.g.

$$\vec{E} = \hat{x} e^{ikx}$$

satisfies $(\nabla^2 + k^2)\hat{x} e^{ikx} = 0$ but not Gauss's law:

$$\begin{aligned} \nabla \cdot \epsilon \vec{E} &= \epsilon \nabla \cdot \hat{x} e^{ikx} \\ &= \epsilon ik e^{ikx} \neq 0 \end{aligned}$$

So we have to be a bit more clever...

To find a Green's function for Maxwell's equation, we will find a vector field which also satisfies the wave equation, but from which a valid \vec{E} satisfying Gauss's law can be obtained.

We will need two theorems:

If $\nabla \cdot \vec{B} = 0$, then there exists a vector field \vec{A} such that $\vec{B} = \nabla \times \vec{A}$.

If $\nabla \times \vec{C} = 0$, then there is a function such that $\nabla f = \vec{C}$.

(for proofs, see a vector calculus text)

From Gauss's law $\nabla \cdot \vec{B} = 0$, by the first theorem there exists \vec{A} such that

$$\vec{B} = \nabla \times \vec{A}$$

\vec{A} is the magnetic vector potential.

From Faraday's law,

$$\nabla \times \vec{E} = i\omega \vec{B} = i\omega \nabla \times \vec{A}$$

or

$$\nabla \times (\vec{E} - i\omega \vec{A}) = 0$$

By the second theorem, there exists ϕ such that

$$\vec{E} - i\omega \vec{A} = -\nabla \phi$$

ϕ is the electric potential. Using this result in Ampere's law gives

$$\nabla \times \left(\frac{1}{\mu} \nabla \times \vec{A} \right) = -i\omega \epsilon \left(-\nabla \phi + i\omega \vec{A} \right) + \vec{J}$$

$$\nabla \times \nabla \times \vec{A} = i\omega \epsilon \mu \nabla \phi + \omega^2 \epsilon \mu \vec{A} + \mu \vec{J}$$

$$-\nabla \times \nabla \times \vec{A} + k^2 \vec{A} = -i\omega \epsilon \mu \nabla \phi - \mu \vec{J}$$

Using $\nabla^2 \vec{A} = -\nabla \times \nabla \times \vec{A} + \nabla \nabla \cdot \vec{A}$,

$$(\nabla^2 + k^2) \vec{A} = -i\omega \epsilon \mu \nabla \phi + \nabla \nabla \cdot \vec{A} - \mu \vec{J}$$

$$(\nabla^2 + k^2) \vec{A} = \nabla(\nabla \cdot \vec{A} - i\omega\epsilon\mu\phi) - \mu \vec{J}$$

Now, we use the fact that

$$\begin{aligned} \vec{B} &= \nabla \times \vec{A} \\ &= \nabla \times (\vec{A} + \nabla\psi) \\ &= \nabla \times \vec{A} + \nabla \times \nabla\psi \\ &= \nabla \times \vec{A} = \vec{B} \end{aligned}$$

to add a gradient term to \vec{A} such that

$$\nabla \cdot \vec{A} = i\omega\epsilon\mu\phi$$

This is known as a "gauge". There are other gauges, or specifications for the divergence of \vec{A} . This one is called the Lorenz gauge.

Using this gauge relationship, we have

$$\boxed{(\nabla^2 + k^2) \vec{A} = -\mu \vec{J}}$$

Now we can use the scalar Green's function for each component:

$$\begin{aligned} \vec{A} &= \mu \iiint d\vec{r}' g(\vec{r}, \vec{r}') \vec{J}(\vec{r}') \\ &= \mu \iiint d\vec{r}' \frac{e^{ik|\vec{r}-\vec{r}'|}}{4\pi|\vec{r}-\vec{r}'|} \vec{J}(\vec{r}') \end{aligned}$$

The electric field is

$$\begin{aligned} \vec{E} &= i\omega\vec{A} - \nabla\phi \\ &= i\omega\vec{A} - \frac{\nabla \nabla \cdot \vec{A}}{i\omega\epsilon\mu} \\ &= i\omega\mu \iiint d\vec{r}' g(\vec{r}, \vec{r}') \vec{J}(\vec{r}') - \frac{1}{i\omega\epsilon\mu} \nabla \nabla \cdot \iiint d\vec{r}' g(\vec{r}, \vec{r}') \vec{J}(\vec{r}') \\ &= i\omega\mu \left[1 + \frac{1}{k^2} \nabla \nabla \cdot \right] \iiint d\vec{r}' g(\vec{r}, \vec{r}') \vec{J}(\vec{r}') \end{aligned}$$

This term cancels out the 'non-Gaussian' solutions to the wave equation!