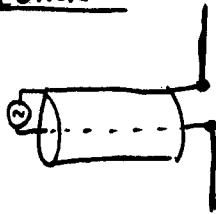


An antenna transforms energy from a transmission line or a waveguide to a wave propagating in free space.

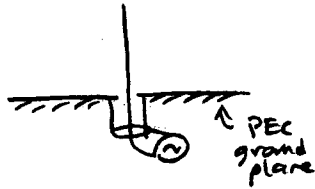
Examples:

Current

Dipole:



Monopole:



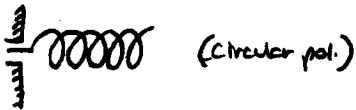
Loop:



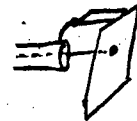
Patch:



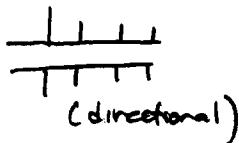
Helical:



"Planar inverted F":
(cell phone)



Yagi:

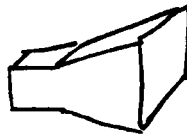


Aperture

Open-ended waveguide:

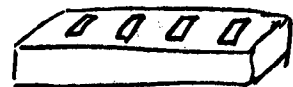


Horn:



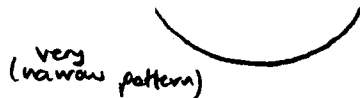
(narrower pattern)

Slotted waveguide:



Feed

Other types: Reflector:



very narrow pattern

Array:

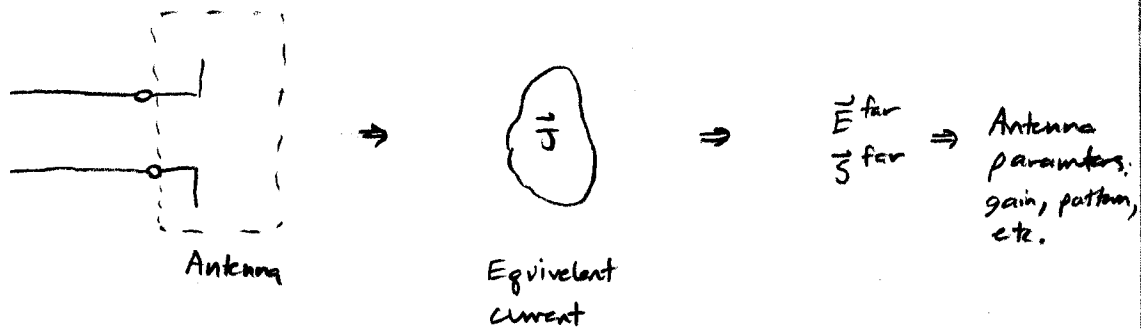


↑ ↑
Ang of the above antennas

(shapeable beam, steerable...)

22-141 50 SHEETS
22-142 100 SHEETS
22-144 200 SHEETS



Goal of antenna analysis:

1. Find an equivalent current.

Approaches: Approximations

- Linear antennas: Hertzian dipole model, triangular current model, sinusoidal, ...
- Aperture antennas: aperture field \approx incident field
- etc.

Numerical method (ECEn 563)

2. Find far fields using radiation integral

3. Find antenna parameters from $\vec{J} \text{ far}$

Antenna design = inverse problem: given antenna parameters, design the antenna. (More difficult).

Basic antenna parameters are:

Radiation pattern, gain, directivity
main lobe beamwidth, sidelobe levels

Radiation resistance

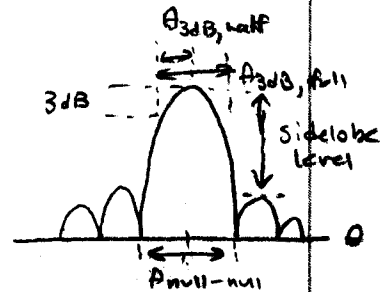
Efficiency

Bandwidth, frequency of operation, size

Radiation pattern

$$\vec{S} = \vec{E} \times \vec{H}^* \rightarrow \underline{f(\theta, \phi)} \frac{1}{r^2} \hat{r}, \quad r \rightarrow \infty$$

→ radiation pattern = $\frac{f(\theta, \phi)}{f_{\max}}$



Gain/directivity

Power radiated to far field: $P_{\text{rad}} = \oint_S \vec{S}_{\text{av}} \cdot d\vec{S}$



Directivity: $D(\theta, \phi) = \frac{\text{Power density at } (\theta, \phi)}{\text{Power density of isotropic radiator}}$

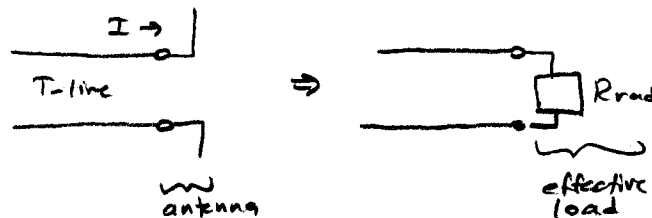
$$= \frac{S_{\text{av}, r}^{\text{ff}}}{P_{\text{rad}} / 4\pi r^2}, \quad S_{\text{av}, r}^{\text{ff}} = \frac{1}{2} \text{Re}[\vec{E} \times \vec{H}^*] \cdot \hat{r}, \quad (\text{large } r \dots \text{ for field})$$

Efficiency: $\eta_{\text{rad}} = \frac{P_{\text{rad}}}{P_{\text{in}}}$, $\eta_{\text{rad}} < 1$ if antenna is lossy

Gain: $G(\theta, \phi) = \eta_{\text{rad}} D(\theta, \phi)$

Radiation resistance

$$R_{\text{rad}} = \frac{P_{\text{rad}}}{\frac{1}{2} |I|^2}$$



Bandwidth/frequency/size

Resonant antenna: desirable radiation resistance over some band

Low frequencies: $R_{\text{rad}} \sim (kR)^2$, so antenna must be large or else efficiency is low

Antenna $\ll \lambda$: broad pattern, low gain

Antenna $\gg \lambda$: narrow beam, high gain - 3 m dish at 1.5 GHz ≈ 30 dB gain

Example: Hertzian dipole

For the Hertzian dipole,

$$\begin{aligned} \vec{S} &= \frac{1}{2} \operatorname{Re} \vec{E} \times \vec{H}^* \\ &= -i\omega\mu \frac{e^{i\mathbf{k}\cdot\mathbf{r}}}{4\pi r} \hat{\theta} I\ell \sin\theta \times \left(-\frac{i\omega\mu}{\eta} \frac{e^{+i\mathbf{k}\cdot\mathbf{r}}}{4\pi r} I\ell \sin\theta \hat{\phi} \right)^* \\ &= \frac{(\omega\mu)^2 (I\ell)^2}{\eta (4\pi r)^2} \sin^2\theta \hat{r} \\ &= \frac{\eta}{2} \left(\frac{kI\ell}{4\pi r} \right)^2 \sin^2\theta \hat{r} \end{aligned}$$

The total power radiated is

$$\begin{aligned} P_{\text{rad}} &= \frac{1}{2} \int_0^{2\pi} \int_0^\pi r^2 \sin\theta \, d\theta \, d\phi \, \eta \left(\frac{kI\ell}{4\pi r} \right)^2 \sin^2\theta \\ &= \eta \left(\frac{kI\ell}{4\pi} \right)^2 2\pi \int_0^\pi \sin^3\theta \, d\theta \\ &= \eta \left(\frac{kI\ell}{4\pi} \right)^2 \cdot \frac{8\pi}{3} \\ &= \eta \frac{(kI\ell)^2}{12\pi} \end{aligned}$$

The gain pattern is

$$\begin{aligned} G(\theta, \phi) &= \frac{S(\theta, \phi)}{P_{\text{rad}}/4\pi r^2} \\ &= \frac{4\pi r^2 \left(\eta \left(\frac{kI\ell}{4\pi r} \right)^2 \sin^2\theta \right)}{\eta \left(\frac{kI\ell}{4\pi} \right)^2 \cdot \frac{8\pi}{3}} \\ &= \left(\frac{4\pi}{\frac{8\pi}{3}} \right) \cdot \sin^2\theta \\ &= \frac{3}{2} \sin^2\theta \end{aligned}$$

where we assume no loss, $\eta_r = 1$.

The radiation resistance is

$$\begin{aligned}
 R_r &= \frac{P_{\text{rad}}}{\frac{1}{2} I_0^2} \\
 &= \eta \frac{(kI_0l)^2}{6\pi \cdot \frac{1}{2} I_0^2} \\
 &= \eta \frac{(kl)^2}{6\pi} \\
 &\approx 120\pi \frac{(kl)^2}{6\pi} \\
 &= \underline{\underline{20(kl)^2}}
 \end{aligned}$$

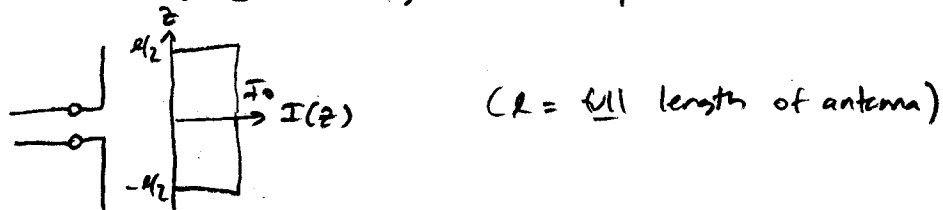
This result holds only for small kl . The reason for this is that if we drive a long wire antenna, the current is not uniform along the wire, so the Hertzian dipole model breaks down.

Hertzian dipole

The Hertzian dipole is

$$\begin{array}{c} \uparrow z \\ \uparrow \oplus q e^{i\omega t} \\ \downarrow \ominus -q e^{-i\omega t} \end{array} \Rightarrow \vec{J} = \delta(x) \delta(y) \delta(z) \hat{z} I \ell$$

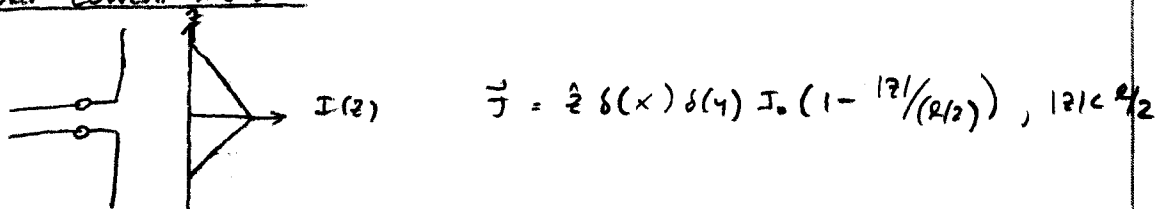
As a linear or wire antenna, this corresponds to the model



where the wires are short ($l \ll \lambda$) and the current on the wires is constant. This leads to

$$R_{\text{rad}} \approx 20 (kl)^2 \approx 200 \Omega \text{ for half-wave dipole, } l = \lambda/2$$

But in reality the current $I(z)$ must vanish at the ends of the wire, so a better current model would be more accurate.

Triangular current model

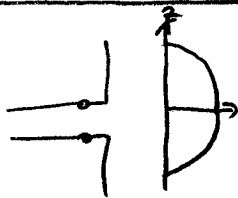
For this current model,

$$\begin{aligned} \vec{F}_i &= \int d\vec{r}' \underbrace{e^{-i\vec{k}\cdot\vec{r}'}}_{\rightarrow 1 \text{ for } l \ll \lambda} \vec{J}(\vec{r}') \hat{z} \\ &= \int_{-l/2}^{l/2} dz \left(1 - \frac{|z|}{l/2}\right) I_0 \hat{z} \\ &= \frac{I_0 \hat{z}}{2} = \frac{1}{2} (\vec{F}_{\text{HD}}) \end{aligned}$$

Because of the $1/2$ factor, R_{rad} decreases by $1/4$:

$$\begin{aligned} R_{\text{rad}} &\approx 5 (kl)^2 \quad (\text{Triangle current model}) \\ &\approx 50 \Omega \text{ for half-wave dipole} \end{aligned}$$

Sinusoidal current model

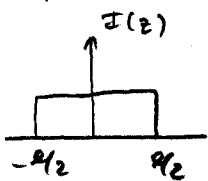


$$\vec{J} = \hat{z} \delta(x) \delta(y) I_0 \sin[k(\ell/2 - |z|)]$$

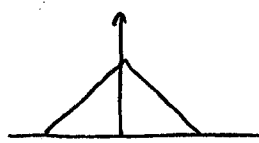
Evaluating \vec{J} for and finding the radiation resistance is

$$R_{rad} \approx 80 \Omega \text{ for half-wave dipole}$$

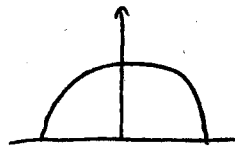
What is the "true" value of R_{rad} ?



Hertzian dipole

200 Ω 

Triangular model

50 Ω 

Sinusoidal

80 Ω (w/ phase, 73 Ω)

true current

Other improvements to the model include finite wire thickness, balun, etc...